

Lecture 24: Ideal Gas Model in Boltzmann Statistics

Lecturer: Qiang Zhu

Scribes: scribe-name1,2,3

24.1 Partition Functions for Composite Systems

Consider a system of just two particles, 1 and 2. If these particles do not interact with each other, their energy is just simply $E_1 + E_2$, then

$$Z_{\text{total}} = \sum_s e^{-\beta[E_1(s)+E_2(s)]} = \sum_s e^{-\beta E_1(s)} e^{-\beta E_2(s)} \quad (24.1)$$

Where the sum runs over all states s for the composite system. If the two particles are *distinguishable*,

$$Z_{\text{total}} = \sum_{s1} \sum_{s2} e^{-\beta E_1(s1)} e^{-\beta E_2(s2)} \quad (24.2)$$

In this case, we split the total partition functions Z into separate Z_1 and Z_2 .

$$Z_{\text{total}} = Z_1 Z_2 \quad (24.3)$$

If the two particles are indistinguishable, we have to apply $1/2$ to reduce double counts.

$$Z_{\text{total}} = \frac{1}{2} Z_1 Z_2 \quad (24.4)$$

This formula is not precisely correct, because there are some terms **in the double sum in which both particles are in the same state when the system is very dense**.

Therefore, the generalization of the equation is the following

$$Z_{\text{total}} = Z_1 Z_2 Z_3 \cdots Z_N \quad (24.5)$$

$$Z_{\text{total}} = \frac{1}{N!} Z_1^N \quad (24.6)$$

24.2 The Partition Function of an Ideal Gas

An ideal gas should have the following partition function form,

$$Z_{\text{total}} = \frac{1}{N!} Z_1^N \quad (24.7)$$

To calculate Z_1 , we must make the Boltzmann factor,

$$e^{-E(s)/kT} = e^{-E_{\text{tr}}(s)/kT} e^{-E_{\text{int}}(s)/kT} \quad (24.8)$$

Since Z is additive,

$$Z_1 = Z_{\text{tr}} Z_{\text{int}} \quad (24.9)$$

Where

$$Z_{\text{tr}} = \sum e^{-E_{\text{tr}}/kT} \quad Z_{\text{int}} = \sum e^{-E_{\text{int}}/kT} \quad (24.10)$$

To calculate Z_{tr} , we can start with the case of a molecule confined to a one-dimensional box.

$$\lambda_n = \frac{2L}{n}, \quad n = 1, 2, \dots, \quad (24.11)$$

$$p_n = \frac{h}{\lambda_n} = \frac{hn}{2L} \quad n = 1, 2, \dots, \quad (24.12)$$

$$E_n = \frac{p_n^2}{2m} = \frac{h^2 n^2}{8mL^2} \quad (24.13)$$

Therefore,

$$Z_{1d} = \sum_n e^{-E_n/kT} = \sum_n e^{-\frac{h^2 n^2}{8mL^2 kT}} \quad (24.14)$$

By doing integration

$$Z_{1d} = \int_0^\infty e^{-\frac{h^2 n^2}{8mL^2 kT}} dn = \frac{\sqrt{\pi}}{2} \sqrt{\frac{8mL^2 kT}{h^2}} = \sqrt{\frac{2\pi m kT}{h^2}} L = \frac{L}{L_Q} \quad (L_Q : \text{Quantum length}) \quad (24.15)$$

For 3 dimension,

$$E_{\text{tr}} = \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m} \quad (24.16)$$

$$Z_{\text{tr}} = \frac{L_x}{L_Q} \frac{L_y}{L_Q} \frac{L_z}{L_Q} = \frac{V}{V_Q} \quad (V_Q : \text{Quantum Volume}) \quad (24.17)$$

$$Z_1 = \frac{V}{V_Q} Z_{\text{int}} \quad (24.18)$$

where

$$Z = \frac{1}{N!} \left(\frac{V Z_{\text{int}}}{V_Q} \right)^N \quad (24.19)$$

and

$$\ln Z = N[\ln V + \ln Z_{\text{int}} - \ln N - \ln V_Q + 1] \quad (24.20)$$

Knowing Z can help us to compute many quantities,

$$U = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{\partial \ln Z}{\partial \beta} \quad (24.21)$$

$$U = -N \frac{\partial \ln Z_{\text{int}}}{\partial \beta} + N \frac{1}{V_Q} \frac{\partial V_Q}{\partial \beta} = N \bar{E}_{\text{int}} + N \frac{3}{2\beta} = U_{\text{int}} + \frac{3}{2} NkT \quad (24.22)$$

$$C_V = \frac{\partial U}{\partial T} = \frac{\partial U_{\text{int}}}{\partial T} + \frac{3}{2} Nk \quad (24.23)$$

$$F = -kT \ln Z = -NkT [\ln V - \ln N - \ln V_Q + 1] + F_{\text{int}} \quad (24.24)$$

From this, it is easy to compute

$$P = -\left(\frac{\partial F}{\partial V}\right)_{T,N} = \frac{NkT}{V} \quad (24.25)$$