

## Lecture 10: Entropy and Pressure

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### 10.1 Mechanical Equilibrium and Pressure

In the last lecture, we just learned the relation between  $S$  and  $T$ , is there any analogy between  $S$  and  $P$ ?

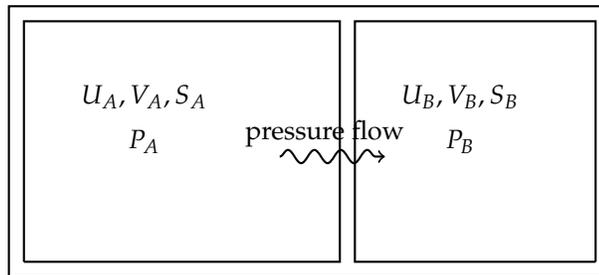


Figure 10.1: A schematic pressure flow between two gases.

Again, we start from the condition when the system reaches its equilibrium,

$$\frac{\partial S_{\text{total}}}{\partial U_A} = 0 \quad \rightarrow \quad \frac{\partial S_{\text{total}}}{\partial V_A} = 0 \quad (10.1)$$

as  $S$  is a function of  $U$  and  $V$ .

We already applied the 1st condition in the previous lecture. How about the 2nd condition?

$$\frac{\partial S_A}{\partial V_A} + \frac{\partial S_B}{\partial V_A} = 0 \quad \rightarrow \quad \frac{\partial S_A}{\partial V_A} = -\frac{\partial S_B}{\partial V_B} \quad (10.2)$$

What's the physical meaning of  $\partial S_A / \partial V_A$ ?

If we dig a bit on the units, we will find  $\partial S_A / \partial V_A$  has a unit of  $\text{N}\cdot\text{m}/\text{K}$ , about  $P/T$  hence we guess

$$\frac{P}{T} = \left(\frac{\partial S}{\partial V}\right)_{U,N} \quad \rightarrow \quad P = T \left(\frac{\partial S}{\partial V}\right)_{U,N} \quad (10.3)$$

Recall that we know how to calculate  $S$ ,

$$S = Nk \ln V + 3/2 Nk \ln U - Nk \ln(f(N)) \quad (10.4)$$

$$P = T \left(\frac{\partial S}{\partial V}\right) = \frac{NkT}{V} \quad (10.5)$$

$$PV = NkT \quad (10.6)$$

Again, we proved the ideal gas law.

## 10.2 Thermodynamic Identity

From the above sections, it seems that  $\Delta S$  can be divided into two parts,

1.  $\Delta U$ , to account for the heat flow
2.  $\Delta V$ , to account for the pressure flow

Let's say,

$$\Delta S = \left(\frac{\Delta S}{\Delta U}\right)\Delta U + \left(\frac{\Delta S}{\Delta V}\right)\Delta V \quad (10.7)$$

Suppose each step is very small, we use

$$dS = \left(\frac{\partial S}{\partial U}\right)dU + \left(\frac{\partial S}{\partial V}\right)dV \quad (10.8)$$

$$dS = \frac{dU}{T} + \frac{PdV}{T} \quad (10.9)$$

$$TdS = dU + PdV \quad (10.10)$$

$$dU = TdS - PdV \quad (10.11)$$

This is the **Thermodynamic Identity**. If you compare it with the 1st law, it just substitutes  $TdS$  with  $Q$ , which is actually the old definition of entropy.

1.  $\Delta U = 0, TdS = PdV$
2.  $\Delta V = 0, dU = TdS$

**Exercise** Under constant entropy

$$\left(\frac{\partial S}{\partial U}\right)dU + \left(\frac{\partial S}{\partial V}\right)dV = 0 \quad (10.12)$$

$$dU = -PdV \quad (10.13)$$

isentropic = quasistatic + adiabatic

$$\Delta S = S_f - S_i = \int_{T_i}^{T_f} \frac{C_P}{T} dT \quad (10.14)$$

$$S(300K) = S(0K) + C_P \int_0^{300} \frac{1}{T} dT = 5.8 + 3.5 \cdot 8.31 \cdot \ln(300) = 173.89 \text{ J/K.}$$

This value looks much smaller than the reference value in the appendix (197.67 J/K), because a constant volume assumption is not realistic. A more realistic solution is

$$\Delta S = C_V \ln \frac{P_B}{P_A} + C_P \ln \frac{V_B}{V_A} \quad (10.15)$$

when you consider  $Q = \Delta U - W$ .

$$\Delta S = \frac{Q}{T} \quad \text{quasistatic} \quad (10.16)$$

$$\Delta S > \frac{Q}{T} \quad \text{in practice} \quad (10.17)$$

1. Very fast compression
2. free expansion

### **10.3 Homework**

Problem 3.5, 3.8, 3.11, 3.14, 3.16, 3.27, 3.30, 3.31, 3.32, 3.33